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Properties of aqueous solutions of lithium and calcium chlorides: formulations for use in air conditioning equipment design

Manuel R. Conde

M. Conde Engineering, Eugen-Huber-Str. 61, CH-8048 Zürich, Switzerland Received 4 March 2003; accepted 17 September 2003

Abstract

The dehydration of air, for air conditioning purposes, either for human comfort or for industrial processes, is done most of the times by making it contact a surface at a temperature below its dew point. In this process not only is it necessary to cool that surface continuously, but also the air is cooled beyond the temperature necessary to the process, thus requiring reheating after dehumidification. Although the equipment for this purpose is standard and mostly low-cost, the running costs are high and high grade energy is dissipated at very low efficiency. Alternative sorption-based processes require only low grade energy for regeneration of the sorbent materials, thus incurring lower running costs. On the other hand, sorption technology equipment is usually more expensive than standard mechanical refrigeration equipment, which is essentially due to their too small market share. This paper reports the development of calculation models for the thermophysical properties of aqueous solutions of the chlorides of lithium and calcium, particularly suited for use as desiccants in sorptionbased air conditioning equipment. This development has been undertaken in order to create consistent methods suitable for use in the industrial design of liquid desiccant-based air conditioning equipment. We have reviewed sources of measured data from 1850 onwards, and propose calculation models for the following properties of those aqueous solutions: Solubility boundary, vapour pressure, density, surface tension, dynamic viscosity, thermal conductivity, specific thermal capacity and differential enthalpy of dilution. 2003 Elsevier SAS. All rights reserved.

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1. Introduction

Aqueous solutions of alkali halides have a set of very useful characteristics for applications in refrigeration and air conditioning. While they were first studied and characterized as freezing point depressors and used as brines in refrigeration applications, other of their properties, such as the high water affinity, were later recognized as very interesting for the dehydration of gases, in particular of air. Dehydration (dehumidification) of air is an important operation in many processes, especially in air conditioning, either for human comfort or for industrial processes. The important point here is that the air does not have to be cooled below its dew-point, as is typical of surface dehumidification. It may be sufficiently dehumidified at temperatures well above its dew-point. This may have a significative impact in the energy requirements and on the size and capacity of the cooling equipment necessary.

The R&D effort in this field has increased significatively in the last twenty years, in particular because of the perceived need to replace the more conventional technologies that use mechanical refrigeration equipment. Uncertainties regarding the availability of safe and suitable working fluids for mechanical refrigeration systems, as well as the effort to reduce greenhouse gases emitted in power generation, are pressing manufacturers, designers and users to search for alternative solutions, naturally without any penalty in the perceived quality of life and in the quality of manufacturing facilities.

On the other hand, systems using sorption processes to dehydrate gases require only low grade thermal driving energy for the regeneration of the sorbent materials, besides some electric power for the prime movers in the system, mostly electric motors. Low grade thermal energy for the regeneration process, may be obtained from effluents in many production processes including CHP, or from solar thermal

E-mail address: manuel.conde@mrc-eng.com (M.R. Conde).

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Nomenclature

collectors. Theoretically, at least, such low grade energy should be much cheaper than the electric power required to drive conventional mechanical refrigeration equipment. However, as in many situations in life, the newcomer must overdo the incumbent by a large advantage to be seriously considered.

Basic materials used in sorption based dehydration of air and other gases, are mostly naturally occurring substances. However, besides having to overdo the incumbent technology, sorption based technologies for air conditioning seem to get much less R&D effort than that invested in the development of new, completely artificial substances, suitable to replace those working fluids that had to be banned, due to their hazardous effects on the environment. It is perhaps a signal of our times that a larger effort is invested in the development of non-natural working fluids than on the development of the processes and equipment that may use natural substances to the same effect.

In this communication we report on the properties of the aqueous solutions of lithium and calcium chlorides, particularly on empirical formulations of their thermophysical properties required in the industrial design of sorption based air conditioning equipment. To that end, we have reviewed the literature spanning the period from 1850 onwards. Special care has been taken in order for the formulations to reproduce the measured data sets with good accuracy in their range of validity, and that they remain congruent at the limits of concentration.

Formulations for the calculation of the following properties are described herein:

- Solubility boundary;
- Vapour pressure;
- Density;
- Surface tension;
- Dynamic viscosity;
- Thermal conductivity;
- Specific thermal capacity;
- Differential enthalpy of dilution.

We expect that this paper will contribute to ease the efforts of the few R&D groups and manufacturers active in the development of sorption based technologies for air conditioning applications.

2. Solubility boundary

The solubility boundaries of these two salt solutions are defined by several lines. For salt concentrations lower than that at the eutectic point, the solubility line defines the conditions at which ice crystals start to form. This is the ice line. For higher concentrations, the solubility boundary defines the conditions at which salt hydrates or anhydrous salt crystalize from the solution. This is the crystallization line.

For LiCl–H₂O solutions, the crystallization line, $A-B-$ C–D–E in Fig. 1, defines also the transition points separating the ranges of formation of the various hydrates, as indicated in the figure. Equations for each range were adjusted to experimental data from the literature [1–14] and have the general form for the crystallization line:

$$
\theta = \sum_{i=0}^{2} A_i \xi^i, \qquad \theta \equiv \frac{T}{T_{c,H_2O}}
$$

where *ξ* is the mass fraction of the salt in the solution. For the ice line, the equation is slightly different:

$$
\theta = A_0 + A_1 \xi + A_2 \xi^{2.5}
$$

The parameters *Ai* are included in Table 1, for each range of the boundary.

For $CaCl₂-H₂O$ solutions, the crystallization line is more complex, particularly due to the formation of various tetrahydrates. Some of these are metastable. In the boundary represented in Fig. 2 we decided to show only those generally reported as stable in the literature, i.e., the *α* and *β* tetrahydrates, with the boundaries described below.

Fig. 1. Solubility boundary of aqueous solutions of lithium chloride.

Table 1 Parameters of the equations describing the solubility boundary of LiCl– H2O solutions

Boundary	A ₀	A ₁	A ₂
Ice Line	0.422088	-0.090410	-2.936350
$LiCl-5H2O$	-0.005340	2.015890	-3.114590
$LiCl-3H2O$	-0.560360	4.723080	-5.811050
$LiCl-2H2O$	-0.315220	2.882480	-2.624330
$LiCl-H2O$	-1.312310	6.177670	-5.034790
LiCl	-1.356800	3.448540	0.0

The equations adjusted to measured data from the literature [9,15–25] have the general form given above, with the exception of the ice line, which is described by the equation

$$
\theta = \sum_{i=0}^{2} A_i \xi^i + A_3 \xi^{7.5}
$$

with θ as defined before. The parameters of this equation and those for the various ranges of the crystallization line are given in Table 2.

The transition points between the boundaries are shown in Fig. 2.

Fig. 2. Solubility boundary of aqueous solutions of calcium chloride.

Table 2 Parameters of equations describing the solubility boundary of $CaCl₂–H₂O$ solutions

Boundary	Aο	A1	A2	A_3
Ice Line	0.422088	-0.066933	-0.282395	-355.514247
$CaCl2-6H2O$	-0.378950	3.456900	-3.531310	0.0
CaCl ₂ -4H ₂ O α	-0.519970	3.400970	-2.851290	0.0
CaCl ₂ -4H ₂ O β	-1.149044	5.509111	-4.642544	0.0
$CaCl2-2H2O$	-2.385836	8.084829	-5.303476	0.0
$CaCl2-H2O$	-2.807560	4.678250	0.0	0.0

3. Vapour pressure

The equilibrium pressure of saturated water vapour above aqueous solutions of the salts here considered is perhaps their best studied property. Despite the amount of experimental data, spanning almost a century of research, no consistent formulation for the accurate prediction of this property has yet been published. Most attempts made are limited to short ranges either on concentration or temperature. Formulations derived from basic principles have not done better than the empirical ones and are confined to very dilute solutions. The formulations we report in the following reproduce quite accurately the known experimental data and default to congruent values at the boundaries.

Fig. 3. Relative vapour pressure of aqueous solutions of lithium chloride.

Table 3 Parameters for the vapour pressure equation

	$LiCl-H2O$	$CaCl2-H2O$
π_0	0.28	0.31
π_1	4.30	3.698
π_2	0.60	0.60
π_3	0.21	0.231
π_4	5.10	4.584
π_5	0.49	0.49
π_6	0.362	0.478
π 7	-4.75	-5.20
π_8	-0.40	-0.40
π ^Q	0.03	0.018

We formulate it in terms of the relative vapour pressure, π (relative to water at the same temperature, Appendix A). In fact we established an accurate formula for a single temperature, for which accurate measured data are known, as function of the mass fraction, and derived a correction function for other temperatures in the available range of data.

The general equation is

$$
\pi \equiv \frac{p_{sol}(\xi, T)}{p_{\text{H}_2\text{O}}(T)} = \pi_{25} f(\xi, \theta)
$$

where

$$
f(\xi, \theta) = A + B\theta
$$

$$
A = 2 - \left[1 + \left(\frac{\xi}{\pi_0}\right)^{\pi_1}\right]^{\pi_2}, \qquad B = \left[1 + \left(\frac{\xi}{\pi_3}\right)^{\pi_4}\right]^{\pi_5} - 1
$$

Fig. 4. Relative vapour pressure of aqueous solutions of calcium chloride.

$$
\pi_{25} = 1 - \left[1 + \left(\frac{\xi}{\pi_6}\right)^{\pi_7}\right]^{\pi_8} - \pi_9 e^{-\frac{(\xi - 0.1)^2}{0.005}}
$$

and the parameters are given for both salt solutions in Table 3.

The parameters for $f(\xi, \theta)$ were derived from the data of Gibbard $[26]$ for LiCl–H₂O and from the data of Baker and Waite $[27]$ for CaCl₂-H₂O. The graphs depicted in Figs. 3 and 4, for LiCl–H₂O and CaCl₂–H₂O, respectively, show plots of the vapour pressure along isotherms calculated with the equations adjusted, against the data from the literature. The literature sources are, for LiCl–H₂O [1,2,22,28–42], and for CaCl₂–H₂O [17,18,22,28,30,43–52]. Two types of charts are commonly used to represent the vapour pressure of solutions, namely Dühring [53] and Othmer [54,55] charts. The Dühring chart represents the dew-point temperature of the vapour phase against the bubble point temperature of the solution, while the Othmer chart represents the vapour pressure against the solution temperature. These charts may be easily obtained using the equation above together with that for the water vapour pressure (Appendix A). For the particular application of these solutions as desiccants in air conditioning, a new type of chart showing lines of equilibrium of solution vapour pressure with the partial pressure of water vapour in the air is especially useful. The chart depicted in Fig. 5 for aqueous solutions of lithium chloride, at an atmospheric pressure of 101.325 kPa (sea level atmospheric pressure) as an example, is of this type.

Fig. 5. Chart of equilibrium between the vapour pressure of aqueous solutions of lithium chloride and the partial pressure of water vapour in the air, at the normal sea level atmospheric pressure.

This chart permits the visualization of the required dilution of the solution for a given dehydration of the air, when the efficiency of the dehumidification system is known.

4. Density

The densities of aqueous solutions of lithium and calcium chlorides have been extensively studied and reported upon in the literature (for LiCl–H₂O [1,33,38,56–70] and for CaCl₂– H2O [46,58,62,70–78]).As shown in Figs. 6 and 7, for LiCl– H_2O and $CaCl_2-H_2O$, respectively, the relative (to saturated liquid water at the same temperature) densities may be represented by a single cubic function of the mass fraction ratio solute solvent $\xi/(1 - \xi)$. This function has the form

$$
\rho_{sol}(\xi, T) = \rho_{\rm H_2O}(T) \sum_{i=0}^{3} \rho_i \left(\frac{\xi}{1 - \xi}\right)^i
$$

where the ρ_i are given in Table 4 for the solutions of both chlorides. $\rho_{H_2O}(T)$ is the density of saturated liquid water at the temperature *T* . It is calculated from

$$
\rho_{\text{H}_2\text{O}}(\tau) = \rho_{c,\text{H}_2\text{O}} \big(1 + B_0 \tau^{1/3} + B_1 \tau^{2/3} + B_2 \tau^{5/3} + B_3 \tau^{16/3} + B_4 \tau^{43/3} + B_5 \tau^{110/3} \big)
$$

Fig. 6. Relative densities of aqueous solutions of lithium chloride.

Fig. 7. Relative densities of aqueous solutions of calcium chloride.

Fig. 8. Relative values of the surface tension of aqueous solutions of lithium chloride. Model plotted against literature data.

Table 4 Parameters of the density equation

	$LiCl-H2O$	CaCl ₂ –H ₂ O
ρ_0	1.0	1.0
ρ_1	0.540966	0.836014
ρ_2	-0.303792	-0.436300
ρ_3	0.100791	0.105642

Table 5

Parameters for the liquid water density equation

	Bi
0	1.9937718430
	1.0985211604
っ	-0.5094492996
3	-1.7619124270
	-44.9005480267
	-723692.2618632

where $\tau \equiv 1 - \theta$ and the B_i are given in Table 5. ρ_{c,H_2O} is the density of water at the critical point (322 kg⋅m⁻³).

For solutions of LiCl, the scatter is almost negligible. The break point defined on the base of the data of Applebey et al. [1] for the crystallization boundary, does not exactly match the corresponding transition point in Fig. 1. This is due to the scarce density data for concentrations higher than that at that transition point.

For solutions of $CaCl₂$, the scatter is somewhat larger, although the prediction of a single function remains excellent.

Fig. 9. Relative values of the surface tension of aqueous solutions of calcium chloride. Model plotted against literature data.

The range of application of the equation is $0 \le \xi \le 0.56$ and $0 \le \xi \le 0.60$ for solutions of LiCl and CaCl₂, respectively.

5. Surface tension

The values of the surface tension of aqueous solutions of the lithium and calcium chlorides reported in the literature (for LiCl–H₂O [65,79–86], and for CaCl₂–H₂O [74,87–91]) are well reproduced by a function of the reduced temperature and mass fraction of the form

$$
\sigma_{sol}(\xi,\theta) = \sigma_{H_2O}(\theta) \left(1 + \sigma_1 \xi + \sigma_2 \xi \theta + \sigma_3 \xi \theta^2 + \sigma_4 \xi^2 + \sigma_5 \xi^3\right)
$$

where θ is, as defined before, the reduced temperature of water $\frac{T}{T_{c,H_2O}}$.

The parameters σ_i are given in Table 6 for the solutions of both chlorides. Figs. 8 and 9 show comparisons of values calculated with the model proposed with data from the literature cited above.

$$
\sigma_{\mathrm{H}_2\mathrm{O}}(\theta) = \sigma_0 \big[1 - b(1-\theta) \big] (1-\theta)^{\mu}
$$

The surface tension of water $\sigma_{H_2O}(\theta)$ is calculated with the equation proposed by the IAPWS (International Association for the Properties of Water and Steam) [92], where $\sigma_0 =$ 235.8 mN·m⁻¹, *b* = -0.625 and μ = 1.256.

Fig. 10. Dynamic viscosity of aqueous solutions of lithium chloride. Comparison of model with literature data.

Table 6

Parameters of the equation to calculate the surface tension of aqueous solutions of lithium and calcium chlorides

σ_1	σ	σ 3	σ_A	σ
	LiCl-H ₂ O 2.757115 -12.011299 14.751818 2.443204 -3.147739			
	$CaCl2-H2O$ 2.33067 -10.78779 13.56611 1.95017 -1.77990			

6. Dynamic viscosity

The dynamic viscosity of aqueous solutions of the lithium and calcium chlorides may be calculated with the equation

$$
\eta_{\text{sol}}(\zeta,\theta) = \eta_{\text{H}_2\text{O}}(\theta) e^{\eta_1 \zeta^{3.6} + \eta_2 \zeta + \eta_3 \frac{\zeta}{\Theta} + \eta_4 \zeta^2}
$$

where ζ is defined as $\zeta = \frac{\xi}{(1-\xi)^{1/0.6}}$.

The parameters η_i are those in Table 7 for the solutions of both chlorides. Figs. 10 and 11 show comparisons of calculated values with measured data from the literature. For aqueous LiCl solutions, the data available [67–70,93– 96] agree very well with those calculated, while for the solutions of $CaCl₂$ [70,78,94,95,97,98] some of the older data, in particular the data of Stakelbeck and Plank [109] at -20 °C, show noticeable deviations.

The dynamic viscosity of water $\eta_{\text{H}_2\text{O}}$, considered here at liquid saturation conditions, is calculated from the formulation recommended by the IAPWS [99] for industrial use, for temperatures above 0° C, Appendix B.

Fig. 11. Dynamic viscosity of aqueous solutions of calcium chloride. Comparison of model with literature data.

Table 7 Parameters of viscosity equation for solutions of lithium and calcium chlorides

For subcooled water ($\vartheta < 0$ °C), although the IAPWS formulation would be satisfactory down to approximately -20 °C, we have fitted an equation to the few data available in the literature [100–102], that reproduces those data with excellent accuracy, as shown in Fig. 12. The equation is

$$
\eta_{\text{H}_2\text{O}} = \eta_{\text{H}_2\text{O},0} \times \left(A + B\theta^{0.02} + C\theta^{0.04} + D\theta^{0.08} + E\theta^{2.85} + F\theta^{8} \right)
$$

with *θ* for this particular case defined as

$$
\theta \equiv \frac{T}{228} - 1
$$

 $\eta_{\text{H}_2\text{O},0}$ is the viscosity of liquid water at 0 °C, calculated from the IAPWS formulation. The parameters *A...F* are given in Table 8.

Fig. 12. Dynamic viscosity of subcooled water. Model and literature data.

Table 8 Parameters of the equation for the dynamic viscosity of subcooled water

1.0261862 12481.702 -19510.923 7065.286 -395.561 143922.996			

7. Thermal conductivity

The thermal conductivity of aqueous solutions of lithium and calcium chlorides are, perhaps, their less well-known property. In fact, literature is sparse and the measured data reported seems not to be always beyond doubts regarding accuracy (see for instance, the discussion by Riedel [103] of the results of Meyer [104] and Rau [105] and the answer by Meißner [106]). Even new data, such as those by Assael et al. [107], for calcium chloride and by Takeuchi et al. [108], for lithium chloride, seem to suffer from large uncertainties. This may eventually be due to the very nature of the solutions to be measured and the methods used.

We have considered the data of Riedel [109] and Uemura [96] for lithium chloride, and of Riedel [103,109] for calcium chloride, to obtain the parameters of a model proposed by Riedel [109] and described in the following.

Riedel defined a characteristic value called 'equivalent thermal conductivity depression' and understood it as a value constant for each salt in aqueous solution, at least for diluted solutions:

$$
\alpha_R \equiv \frac{\lambda_{\text{H}_2\text{O}}(T) - \lambda_{\text{sol}}(T, \xi)}{\zeta_{\text{eq}}}
$$

In this equation, *ζ*eq is the 'equivalent ionic concentration' as named by Riedel,

$$
\zeta_{\text{eq}} = \frac{\xi \times \rho_{\text{sol}}(T,\xi) \times I_s}{M}
$$

 I_s is the ionic strength of the species in solution (1 for LiCl, 2 for $CaCl₂$).

Riedel also noted that α_R would tend to decrease at the higher concentrations in his measurements. We found that α_R shows a linear dependence upon concentration, for

Fig. 13. Thermal conductivity of aqueous solutions of lithium chloride. Comparison of model and measured data.

solutions of both lithium and calcium chlorides, decreasing with concentration. Equations for α_R were fitted to the data of Riedel at 20 ◦C.

$$
\alpha_R = \alpha_0 + \alpha_1 \xi
$$

Their parameters are given in Table 9 for both chlorides.

As may be seen in Fig. 13, the data of Uemura [96], for lithium chloride at temperatures from 10 to 90 \degree C, are predicted with good accuracy on this basis.

For calcium chloride the values measured at 20° C are as well predicted with good accuracy, Fig. 14. All other data, as already mentioned, seem to suffer from poor accuracy, and above all from inconsistency. This was discussed by Riedel [103] for the data of Meyer [104] and Rau [105]. On the other hand, Assael et al. [107] data for calcium chloride show an almost linear behaviour with temperature and a relation to water values that does not conform to Riedel's '*equivalent thermal conductivity depression*' model, Fig. 15. Although those last authors suggest that Riedel's model would not explain the dependency of the

Fig. 14. Thermal conductivity of aqueous solutions of calcium chloride. Comparison of model with measured data.

thermal conductivity of the solution upon temperature, we believe that the behaviour shown by their data has essentially to do with their method of measurement and the very nature of the solution to be measured.

 $\lambda_{H_2O}(T)$ is the thermal conductivity of pure liquid water. For temperatures $\vartheta \geqslant 20 \degree C$, it is calculated using the IAPWS formulation for industrial applications as described in Appendix C. For subcooled water we follow the recommendations of Riedel [103] also adopted by McLaughlin [121], which gives a linear variation of the thermal conductivity with temperature:

$$
\lambda_{\rm H_2O}(\theta) = \lambda_{\rm H_2O,20} \times (0.208495 + 1.747278\theta)
$$

 $\lambda_{H_2O,20}$ is the thermal conductivity at 20 °C, calculated with the IAPWS formulation, and θ is the water reduced temperature.

8. Specific thermal capacity

The specific thermal capacities of aqueous solutions of lithium and calcium chlorides were extensively studied due to their earlier use as a brines in refrigeration systems. The range of temperatures considered was that mostly suitable for this application. Measurements have been reported for lithium chloride by Richards and Rowe [110] Jauch [111] Lange and Dürr [112] Gucker and Schminke [113] and

Fig. 15. Thermal conductivity of aqueous solutions of calcium chloride. Temperature dependence.

Bennewitz and Kratz [114]. Rüterjans et al. [115] carried out measurements for temperatures up to 90 ◦C, albeit for low concentrations of the salt, Fig. 18. For calcium chloride the measurements reported by Dickinson et al. [24], Koch [116], Richards and Dole [117] all cover the range of applications as a refrigeration brine. Ruckov [75] extended the range of measurements up to 75 °C, Fig. 19. Tucker [120] reports measurements of the specific thermal capacity for solutions of both chlorides. As may be observed from Fig. 18 for lithium chloride, and Fig. 19 for calcium chloride, those measurements do not fit together with the measurements of other authors.

The equation we propose for the calculation of the specific thermal capacity of the solutions of both chlorides, has the following general form:

$$
Cp_{sol}(T, \xi) = Cp_{\text{H}_2\text{O}}(T) \times (1 - f_1(\xi) \times f_2(T))
$$

The specific thermal capacity of liquid water is calculated with a modified Sato [118] equation fitted to the data of Angell et al. [119] as

$$
Cp_{\text{H}_2\text{O}}(\theta) = A + B\theta^{0.02} + C\theta^{0.04} + D\theta^{0.06} + E\theta^{1.8} + F\theta^{8}
$$

$$
\theta = \frac{T}{228} - 1
$$

Fig. 16 depicts a graph of this equation comparing it with the data and includes its parameters for the two ranges of validity.

A B C DEF G H

Table 10 Parameters of the equations for the specific thermal capacity of the aqueous solutions of lithium and calcium chlorides

Fig. 16. Specific thermal capacity of liquid water for temperatures down to -35 °C.

Fig. 17. Specific thermal capacity of aqueous solutions of lithium chloride. Data measured at constant temperature compared with model.

Fig. 18. Specific thermal capacity of aqueous solutions of lithium chloride. Data measured at constant concentration compared with model.

describes the effects of salt concentration upon the specific thermal capacity for the whole range for the solutions of calcium chloride and the up to $\xi \leq 0.31$ for those of lithium chloride. For larger concentrations, it is linear on mass fraction

$$
f_1(\xi) = D + E\xi
$$

\n $f_2(T)$ is,
\n $f_2(\theta) = F\theta^{0.02} + G\theta^{0.04} + H\theta^{0.06}$
\n $\theta = \frac{T}{228} - 1$

for both chlorides, with the parameters given in Table 10.

Figs. 17 and 18 show comparisons of the model with measured data at constant temperature and concentration, respectively, for aqueous solutions of lithium chloride. Figs. 19 and 20 depict the same comparisons for aqueous solutions of calcium chloride.

Fig. 19. Specific thermal capacity of aqueous solutions of calcium chloride. Data measured at constant temperature compared with model.

9. Differential enthalpy of dilution

In the dehydration of gases (e.g., moist air) with aqueous salt solutions, the water vapour is absorbed by the solution in the conditioner, to be later desorbed in the regenerator. For the salt solutions considered here, the absorption (dilution) is an exothermic process, while the desorption (regeneration) requires the supply of thermal energy to the solution. This thermal energy required (or liberated) is larger than that corresponding to the vaporization (or condensation) of pure water. This difference constitutes the energy of dilution, and when referred to the unit mass of water is named *differential enthalpy of dilution*.

We used data from the literature to establish interpolating equations for the differential enthalpy of dilution of aqueous solutions of lithium chloride [40,112] and calcium chloride [46,124]. The equations have the general form:

$$
\Delta h_d = \Delta h_{d,0} \bigg[1 + \left(\frac{\zeta}{H_1}\right)^{H_2} \bigg]^{H_3}
$$

where ζ is defined from the salt mass fraction as

$$
\zeta = \frac{\xi}{H_4 - \xi}
$$

The reference $\Delta h_{d,0}$ is related to the temperature as

$$
\Delta h_{d,0} = H_5 + H_6 \theta
$$

Fig. 20. Specific thermal capacity of aqueous solutions of calcium chloride. Data measured at constant concentration compared with model.

Table 11

Parameters of the differential enthalpy of dilution equations for solutions of lithium and calcium chlorides

	$LiCl-H2O$	CaCl ₂ –H ₂ O
H_1	0.845	0.855
H ₂	-1.965	-1.965
H_3	-2.265	-2.265
H_4	0.6	0.8
H_5	169.105	-955.690
H ₆	457.850	3011.974

The parameters H_i for these equations are given in Table 11.

The graphs of Figs. 21 and 22 show comparisons of the literature data with calculations with the models described, for aqueous solutions of lithium and calcium chlorides, respectively.

10. Discussion

This paper presents a complete collection of interpolating equations for the most important properties of aqueous solutions of the lithium and calcium chlorides, necessary in the design of absorption and air conditioning equipment, based on sorption processes with these salts. The data considered originate from a detailed research of old and new literature and concerns measurements by methods that evolved with

Fig. 21. Calculated and measured values of the differential enthalpy of dilution for aqueous solutions of lithium chloride.

time and the technologies available. This however, should not mean that newer measurements are more accurate than older ones. We have systematically given more weight to the data that represented best the data in the database for a given property, as is evident from the graphs. This means that although we have represented in the graphs all the data available in the literature, not all have been considered in the fitting process. The selection criteria is essentially based on the accuracy reported for the measurements, though this is not always available. In some cases, e.g., the crystallization line of aqueous solutions of calcium chloride, many experimental observations have been excluded. Some of the measurements of other properties seem to suffer from poor control of the concentration, given the strong hydrophilic characteristics of the solutions (e.g., some measurements of the surface tension). It also seems to us that the electrical conductivity of the solutions may have interfered with some of the most recent measurements of the thermal conductivity.

Although we would not like to give specific confidence limits for the values predicted by the equations proposed, we consider them well within the boundaries accepted in engineering calculations. We believe the equations proposed to represent a very significative contribution to the process of computer-assisted design of absorption and liquid desiccant based air conditioning systems. The equations were programmed into MathCad® calculation sheets that may be

Fig. 22. Calculated and measured values of the differential enthalpy of dilution for aqueous solutions of calcium chloride.

used as a base for process design and the design of parts of equipment.

Appendix A. Vapour pressure of water over the liquid phase

The vapour pressure of ordinary water over the liquid phase is calculated with an equation due to Saul and Wagner [123] as follows. The parameters of this equation are given in Table 12

$$
\ln\left(\frac{p}{p_{c,H_2O}}\right)
$$

=
$$
\frac{A_0\tau + A_1\tau^{1.5} + A_2\tau^3 + A_3\tau^{3.5} + A_4\tau^4 + A_5\tau^{7.5}}{1 - \tau}
$$

$$
\tau = 1 - \frac{T}{T_{c,H_2O}}
$$

Appendix B. The IAPWS formulation for the dynamic viscosity of ordinary water substance for industrial use

The dynamic viscosity of liquid water at temperatures above 0° C is calculated with the IAPWS formulation for industrial use [99] as follows:

$$
\bar{\eta} = \bar{\eta}_0(\overline{T}) \times \bar{\eta}_1(\overline{T}, \overline{\rho}) \times \bar{\eta}_2(\overline{T}, \overline{\rho})
$$

The term

$$
\bar{\eta}_0(\overline{T}) = \overline{T}^{0.5} \left(\sum_{i=0}^3 H_i \overline{T}^{-i} \right)^{-1}
$$

represents the viscosity of steam in the ideal gas limit with the parameters H_i given in Table 13.

The second term is

$$
\bar{\eta}_1(\overline{T}, \overline{\rho}) = \exp \left\{ \overline{\rho} \sum_{i=0}^{5} \sum_{j=0}^{6} G_{i,j} (\overline{T}^{-1} - 1)^i (\overline{\rho} - 1)^j \right\}
$$

with the parameters $G_{i,j}$ as given in Table 14.

The term $\bar{\eta}_2(\bar{T}, \bar{\rho})$ may be taken as unity for industrial applications, since it concerns only a narrow region around the critical point. The variables with the bar above represent reduced values. These are reduced as follows: $T = T/T^*$, $\bar{\rho} = \rho/\rho^*$, $\bar{\eta} = \eta/\eta^*$, with the reference values: $T^* =$ 647.226 K, $\rho^* = 317.763 \text{ kg} \cdot \text{m}^{-3}$, $\eta^* = 55.071 \times 10^{-6} \text{ Pa.s.}$

Appendix C. The IAPWS formulation for the thermal conductivity of ordinary water substance for industrial use

This formulation [122], as the foregoing one, is included here for the sake of completeness. The IAPWS formulation

Table 12 Parameters of the equation for the vapour pres-

sure of water over the liquid phase		
	A_i	
	-7.858230	
	1.839910	
2	-11.781100	
\mathcal{R}	22.670500	
4	-15.939300	
5	1.775160	

Table 13 *Hi* parameters for the water viscosity equation

for industrial use consists of the following interpolating equation

$$
\bar{\lambda} = \bar{\lambda}_0(\overline{T}) \times \bar{\lambda}_1(\bar{\rho}) \times \bar{\lambda}_2(\overline{T}, \bar{\rho})
$$

The first term, $\bar{\lambda}_0(\bar{T})$, represents the thermal conductivity of steam in the ideal-gas limit, and is

$$
\bar{\lambda}_0(\overline{T}) = \overline{T}^{0.5} \bigg(\sum_{i=0}^3 L_{0,i} \overline{T}^i \bigg)
$$

The term $\bar{\lambda}_1(\bar{\rho})$ is defined by

$$
\bar{\lambda}_1(\bar{\rho}) = L_{1,0} + L_{1,1}\bar{\rho} + L_{1,2} \exp\{L_{1,3}(\bar{\rho} + L_{1,4})^2\}
$$

and $\lambda_2(T, \bar{\rho})$ is defined by the equation

$$
\bar{\lambda}_2(\overline{T}, \bar{\rho}) = \left(\frac{L_{2,0}}{\overline{T}^{10}} + L_{2,1}\right) \bar{\rho}^{9/5} \exp\{L_{3,0}(1 - \bar{\rho}^{14/5})\} \n+ L_{2,2} \Lambda_0 \bar{\rho}^{A_1} \exp\left\{\left(\frac{A_1}{1 + A_1}\right) (1 - \bar{\rho}^{1+A_1})\right\} \n+ L_{2,3} \exp\left\{L_{3,1} \bar{T}^{3/2} + \frac{L_{3,2}}{\bar{\rho}^5}\right\}
$$

*Λ*⁰ and *Λ*¹ are functions of

$$
\Delta \overline{T} = |\overline{T} - 1| + L_{3,3}
$$

Defined as

Table 15

$$
\Lambda_0 = \begin{cases} \frac{1}{\Delta \overline{T}} \Longleftarrow \overline{T} \geq 1\\ \frac{L_{3,5}}{\Delta \overline{T}^{3/5}} \Longleftarrow \overline{T} < 1\\ \Lambda_1 = 2 + \frac{L_{3,4}}{\Delta \overline{T}^{3/5}} \end{cases}
$$

The parameters of these equations are given in Table 15. The variables with the bar above represent reduced values. These are reduced as follows: $T = T/T^*$, $\bar{\rho} = \rho/\rho^*$, $\lambda =$ *λ*/λ^{*}, with the reference values: $T^* = 647.26$ K, $ρ^* =$ $317.7 \text{ kg} \cdot \text{m}^{-3}$, $\lambda^* = 1.0 \text{ W} \cdot \text{m}^{-1} \cdot \text{K}^{-1}$.

Parameters for the thermal conductivity equation of normal water

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